



# Description of Global EGAM in the maximum of local frequency during current ramp-up discharges in DIII-D

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# Outline

- GAM / EGAM – brief introduction and relation to energetic particle transport.
- Dispersion Relation of GAM / EGAM.
- Theoretical description of global EGAM.
- Comparison to N=0 oscillation during ramp-up discharge on DIII-D – discharge #159243.

# GAM – Geodesic Acoustic Mode

- GAM are a low frequency electrostatic oscillation, with toroidal  $N=0$ , poloidal  $M=0$  combined with  $M=\pm 1, \pm 2$ .
- First predicted by [*N. Winsor, Phys. Fluids (1968)*].
- The eigenmode can appear with reverse  $q$  profile (close to the center), when the dispersion relation presents a maximum for thermal plasma [*H.L. Berk (2006) / V.P. Lakhin (2014)*].
- The presence of a similar mode excited by Energetic Particles (**E-GAM**), theoretically predicted [*G. Y. Fu, (2008)*] and experimentally observed [*R. Nazikian (2008)*] is possibly a cause of energetic ion losses.

# GAM calculated by NOVA Code

- NOVA calculates the dispersion relation based on ideal MHD equations [*C.Z. Cheng, Phys. Fluids (1986)* ]:

$$\begin{aligned} \omega^2 \rho \frac{|\nabla\psi|^2}{B^2} \xi_s + (\mathbf{B} \cdot \nabla) \frac{|\nabla\psi|^2}{B^2} (\mathbf{B} \cdot \nabla) \xi_s + \gamma p k_s \nabla \cdot \vec{\xi} &= 0, \\ \left( \frac{\gamma p}{B^2} + 1 \right) \nabla \cdot \vec{\xi} + \frac{\gamma p}{\omega^2 \rho} (\mathbf{B} \cdot \nabla) \frac{\mathbf{B} \cdot \nabla}{B^2} \nabla \cdot \vec{\xi} + k_s \xi_s &= 0. \end{aligned} \quad (1)$$

where  $\psi$  is the poloidal magnetic flux,  $p$  the plasma pressure,  $\rho$  density,  $\xi_s \equiv \frac{\vec{\xi}[\mathbf{B} \times \nabla\psi]}{|\nabla\psi|^2}$ ,  $\vec{\xi}$  the plasma displacement,  $k_s \equiv \frac{2\mathbf{k} \cdot [\mathbf{B} \times \nabla\psi]}{|\nabla\psi|^2}$ ,  $\mathbf{k}$  and  $\mathbf{B}$  are vector of magnetic curvature and field.

- The adiabatic index  $\gamma = 5/3$ , is a free parameter in the code, was changed to make the MHD [*N. Winsor, Phys. Fluids (1968)*] equal to kinetic result  $\omega_{G,mhd}^2 = \frac{\gamma p}{\rho R^2} \left( 2 + \frac{1}{q^2} \right) = \omega_{G,th}^2 = \left[ \frac{7}{4} + \tau_e + \frac{23+4\tau_e(4+\tau_e)}{2q^2(7+4\tau_e)} \frac{1}{q^2} \right] \frac{v_{Ti}^2}{R^2}$ .
- NOVA calculates the N=0 dispersion relation with slightly lower frequency than the standard formula.
- GAM eigenmode can be calculated for strongly reverse q profile, when  $\omega_G^2$  produces maximum OFF-AXIS (accumulation point).

# Kinetic equation and distribution function

Using the drift kinetic equation in the form:

$$\frac{\partial}{\partial t} f_\alpha(r, \vartheta, t) + \frac{v_{||}}{qR_0} \frac{\partial}{\partial \vartheta} f_\alpha(r, \vartheta, t) - v_{dr} \frac{\partial}{\partial r} f_\alpha(r, \vartheta, t) = e_\alpha \frac{\partial}{\partial E} F_\alpha(r, \vartheta) \left( \frac{v_{||}}{qR_0} \frac{\partial}{\partial \vartheta} \phi(r, \vartheta, t) - v_{dr} \frac{\partial}{\partial r} \phi(r, \vartheta, t) \right) \quad (2)$$

Radial drift velocity  $v_{dr,\alpha} = \sin\vartheta \left( v_{||}^2 + \frac{1}{2} v_\perp^2 \right) / \omega_{c\alpha} R_0$

Maxwellian equilibrium distribution function for thermal particles (*th*) and slowing down like distribution function for (*ep*).

$$F_{th,\alpha} = \frac{n_{0\alpha}(r)}{\pi^{3/2}} \frac{e^{-\left(\frac{(v_{||}^2+v_\perp^2)}{v_T^2}\right)}}{v_T^3}, \quad F_{ep} = \frac{n_{0b}(r)\sqrt{1-\lambda_0}}{\ln(U_b/U_c)\pi^{3/2}} \frac{e^{-\left(\frac{(\lambda-\lambda_0)^2}{\Delta\lambda_0^2}\right)}}{\Delta\lambda_0 U^3}$$

The perturbed frequency  $f_\alpha$  and electrostatic potential  $\phi$  expanded in Fourier for  $\vartheta$ , and eikoidal approximation for  $r$  and  $t$ :

$$X(r, \vartheta, t) = (X_0 + X_s \sin\vartheta + X_c \cos\vartheta + X_{2s} \sin 2\vartheta + X_{2c} \cos 2\vartheta) \exp(ik_r r - i\omega t) \quad (3)$$

Eikoidal approximation imply  $\partial/\partial r \rightarrow ik_r$

# Density Oscillation

- From quasineutrality, in the limit  $\rho^2 k_r^2 \ll 1$
- Electron density oscillations:

$$\begin{aligned} n_{es} &= -\frac{2e_i n_0}{m_i v_{Ti}^2 \tau_e} \phi_s \\ n_{e2c} &= -\frac{2e_i n_0}{m_i v_{Ti}^2 \tau_e} \phi_{2c} \\ n_{ec} &= n_{e2s} = 0 \end{aligned} \tag{4}$$

- Electrostatic oscillations:

$$\begin{aligned} \phi_s &= -\tau_e \left[ \left( \frac{v_{Ti}}{\Omega} + (2 + \tau_e) \frac{v_{Ti}^3}{\Omega^3} \right) q \rho_{th} k_r + (27 + 14\tau_e + 2\tau_e^2) \frac{v_{Ti}^3}{\Omega^3} q^3 \rho_{th}^3 k_r^3 \right] \phi_0 \\ \phi_{2c} &= -\tau_e \left[ (7 + 4\tau_e) \frac{v_{Ti}^2}{\Omega^2} + (161 + 92\tau_e + 20\tau_e^2) \frac{v_{Ti}^4}{16\Omega^4} \right] q^2 \rho_{th}^2 k_r^2 \phi_0 \\ \phi_c &= \phi_{2s} = 0 \end{aligned} \tag{5}$$

- EP effects are negligible due to small density.

# Thermal GAM dispersion relation

- From average of radial current  $\langle j_\alpha \rangle = e_\alpha \oint d^3v v_{dr,\alpha} \sin \vartheta f_\alpha$ .
- The resonance condition  $\langle j_i \rangle + \langle j_e \rangle = 0$ , disregarding  $\rho^2 k_r^2$ .
- Thermal GAM dispersion relation ( $D_{th} = 0$ ):

$$D_{th} = 1 - \frac{7+4\tau_e}{4} \frac{v_{Ti}^2 \sigma_s^2}{\omega^2 R_0^2} + \frac{23+16\tau_e+4\tau_e^2}{8q^2} \frac{v_{Ti}^4 \sigma_s^4}{\omega^4 R_0^4 q^2} + iq^2 2\sqrt{\pi} \left[ \frac{\omega^3 R_0^3 q^3}{v_{Ti}^3} + (1 + \tau_e) \frac{\omega R_0 q}{v_{Ti}} \right] \exp \left( -\frac{\omega^2 R_0^2 q^2}{v_{Ti}^2} \right). \quad (6)$$

$\sigma_s$  is formally included, calculated by NOVA.

- From  $\omega \rightarrow \omega + i\gamma$  with  $\gamma \ll \omega$ , the linear growth rate:

$$\gamma \approx -\frac{\omega R_0 q}{v_{Ti}} q^2 \sqrt{\pi} \left[ \frac{\omega^2 R_0^2 q^2}{v_{Ti}^2} + (1 + \tau_e) \right] \exp \left( -\frac{\omega^2 R_0^2 q^2}{v_{Ti}^2} \right)$$

# EGAM dispersion relation

- EGAM Dispersion relation contribution ( $\langle j_i \rangle + \langle j_e \rangle + \langle j_{ep} \rangle = 0 \rightarrow D_{th} + D_{ep} = 0$ ):

$$D_{ep} = -\frac{n_b}{n_i 16(1-\lambda_0)^2 \ln\left(\frac{V_b}{V}\right)} \left[ \left( (5\lambda_0 - 2)(2 - \lambda_0) - \Delta\lambda_0^2 \frac{3(20 - \lambda_0^2 - 4\lambda_0)}{16(1-\lambda_0)} \right) \ln\left(1 - \frac{\omega_{tr}^2}{\omega^2}\right) + \frac{2\lambda_0(\lambda_0 - 2)\omega_{tr}^2}{\left(1 - \frac{\omega_{tr}^2}{\omega^2}\right)\omega^2} - \Delta\lambda_0^2 \frac{3(2 - \lambda_0)\omega_{tr}^2}{(1-\lambda_0)\left(1 - \frac{\omega_{tr}^2}{\omega^2}\right)^2 \omega^2} \right] \quad (7)$$

- EP toroidal transit frequency  $\omega_{tr} = V_b \sqrt{(1 - \lambda_0)} / qR_0$ .
- The injection angle  $\lambda_0 = V_{b\perp}^2 / V_b^2$ .
- The linear growth rate  $\gamma \propto (5 - 2\lambda_0)(2 - \lambda_0) - \Delta\lambda_0^2 \frac{15}{4(1-\lambda_0)}$ , for  $\omega^2 < \omega_{tr}^2$ .
- **Instability found when  $\lambda_0 > 0.43$**  for  $\Delta\lambda_0 \sim 0.2$ .
- Result similar to [Z. Qiu, (2010)] for  $\Delta\lambda_0 = 0$ ,  $\sigma_s = 1$ .

# Eigenmode Equation via Finite Larmor Radius (FLR)

- Keeping terms up to  $\rho^2 k_r^2$ :

$$\langle j_i \rangle + \langle j_e \rangle + \langle j_{ep} \rangle = 0 \quad \rightarrow \quad [(D_{th} + D_{ep}) + (L_{th}\rho_{th}^2 + L_{ep}\rho_{ep}^2)k_r^2]k_r\phi_0 = 0. \quad (8)$$

- Thermal FLR term:

$$L_{th} = \frac{3}{8} + \frac{747+481\tau_e+140\tau_e^2+\tau_e^3}{64} \frac{v_{Ti}^4}{\omega^4 R_0^4}. \quad (9)$$

- EP FLR-EP term:

$$L_{ep} = \frac{n_b(2-\lambda_0)^3}{n_i 16(1-\lambda_0)^2 \ln(\frac{V_b}{V_c})} \left[ \left( \frac{(9\lambda_0-2)}{1 - \frac{\omega_{tr}^2}{\omega^2}} + \Delta\lambda_0^2 \frac{3(\lambda_0^2+12\lambda_0-76)}{16(1-\lambda_0)(2-\lambda_0)} \right) \left( \ln \left( 1 - \frac{\omega_{tr}^2}{\omega^2} \right) + \frac{1}{16} \ln \left( 1 - 4 \frac{\omega_{tr}^2}{\omega^2} \right) \right) \right. \\ \left. - \frac{3(8\lambda_0^2+20\lambda_0-8)\omega_{tr}^4 - 5\omega_{tr}^2\omega^2(9\lambda_0-2)}{16(\omega^2-\omega_{tr}^2)(\omega^2-4\omega_{tr}^2)} - \Delta\lambda_0^2 \frac{600\omega_{tr}^6\omega^2 - 611\omega_{tr}^4\omega^4 - 112\omega_{tr}^4 + 190\omega_{tr}^2\omega^6}{16(2-\lambda_0)(\omega^2-\omega_{tr}^2)^2(\omega^2-4\omega_{tr}^2)^2} \right] \quad (10)$$

- FLR connect different magnetic surfaces
- In presence of EP, FLR is dominated by  $L_{ep}$ ,  $\rho_{ep}^2 \gg \rho_{th}^2$

# Eigenmode Equation

- Recovering the derivative form:  $\partial/\partial r \sim ik_r$  in eq.8 ( $[(D_{th} + D_{ep}) + (L_{th}\rho_{th}^2 + L_{ep}\rho_{ep}^2)k_r^2]k_r\phi_0 = 0$ )

$$Q(r, \omega) \frac{d\phi_0}{dr}(r, \omega) - \frac{\partial^2}{\partial r^2} \frac{d\phi_0}{dr}(r, \omega) = 0 \quad (11)$$
$$Q(r, \omega) = \frac{D_{th} + D_{ep}}{F_{th}\rho_{th}^2 + F_{ep}\rho_{ep}^2}$$

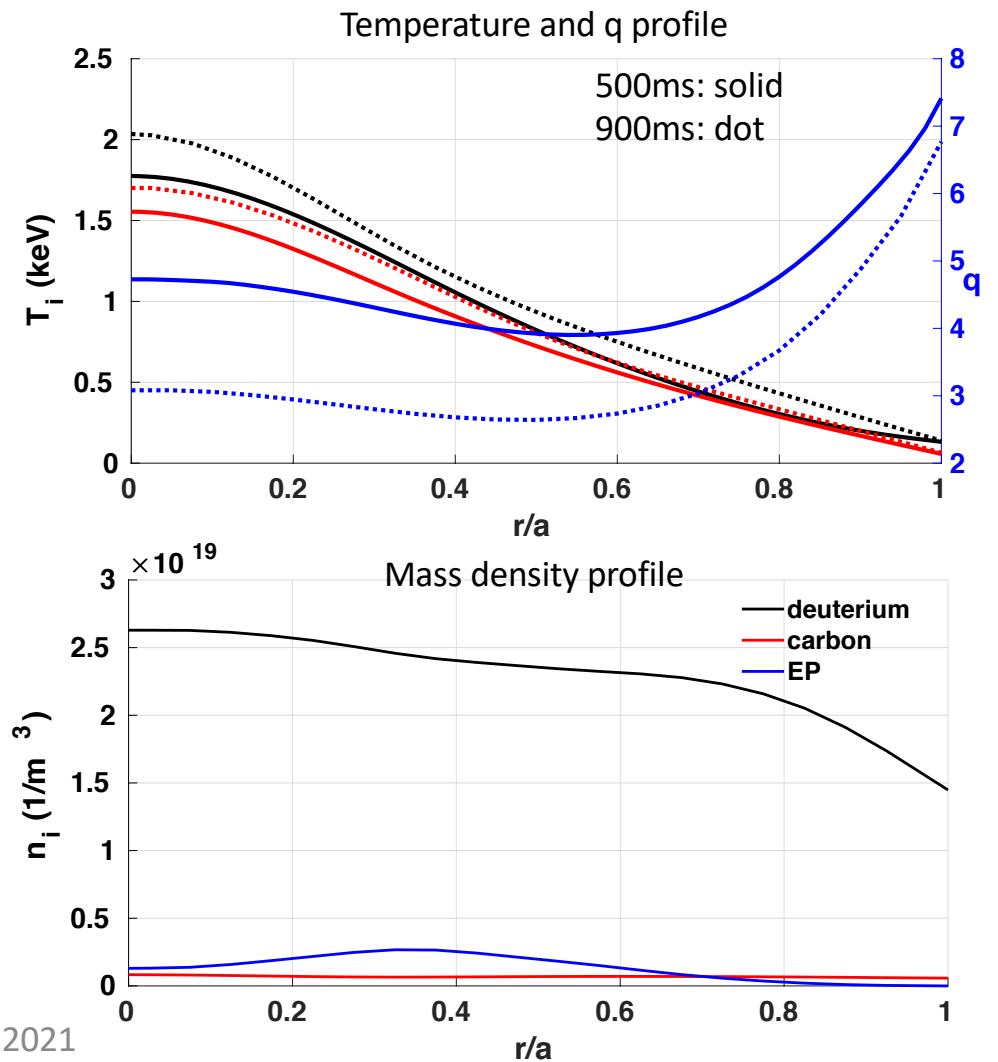
- The solution is calculated numerically between the axis and some point far from  $Q \sim 0$ , then matched with the WKB approximation:

$$\frac{d\phi_0}{dr} = \frac{C_1 e^{i \int \sqrt{Q} dr} + C_2 e^{-i \int \sqrt{Q} dr}}{Q^{1/4}} \quad (12)$$

- A physical solution is found only around the maximum of dispersion relation.
- $Im(Q)$  induces convergence runaway [LeVeque (2007)], physical solution is found for  $C_2=0$ .
- Disregarding  $Im(F_{ep})$  preserves the eigenmode structure, with minimum convergence runaway solution.

# TRANSP profile of #159243

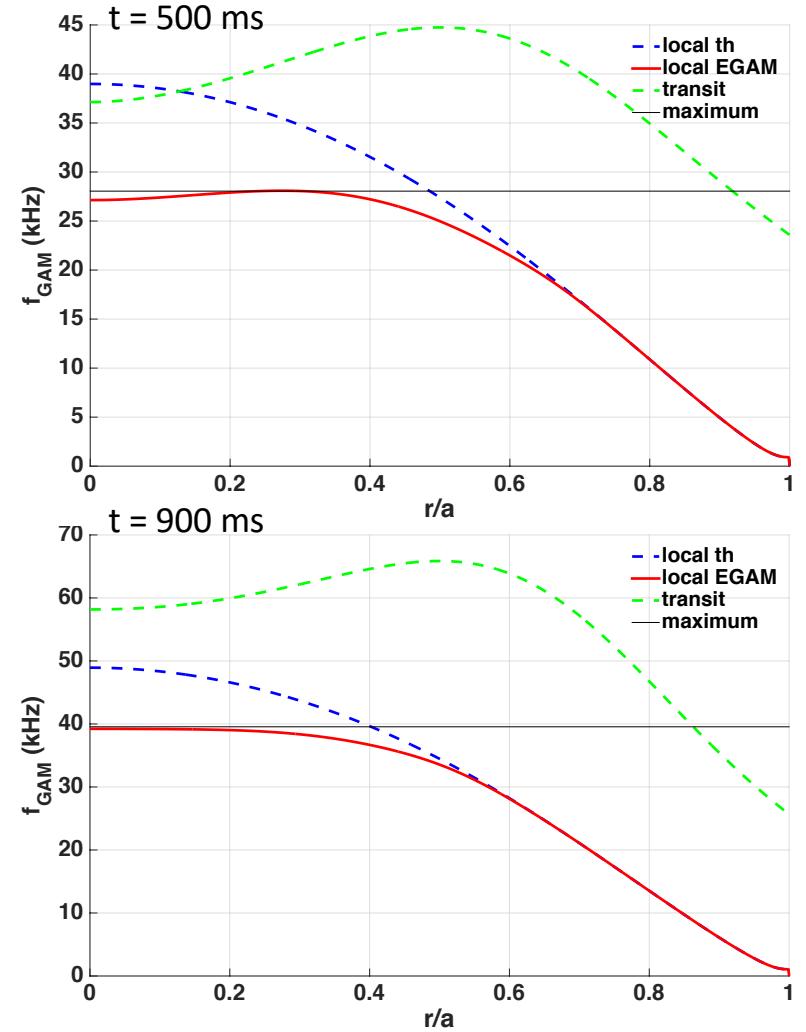
- Time evolution:
  - 500ms – solid line
  - 900ms – dot line
- No strong radial variation in temperature, and the ratio remains around  $\tau_e \sim 0.86$ .
- $q$  at axis decrease from  $\sim 4.7$  to  $3$ .
- mostly co-tangencial discharge with  $\lambda_0 = 0.5\sim0.6$ ,  $\Delta\lambda_0 = 0.2$ .
- Kick model for EP is used [M. Podesta, (2014)].
- Carbon impurities correction are included via effective mass.



# Local dispersion relation

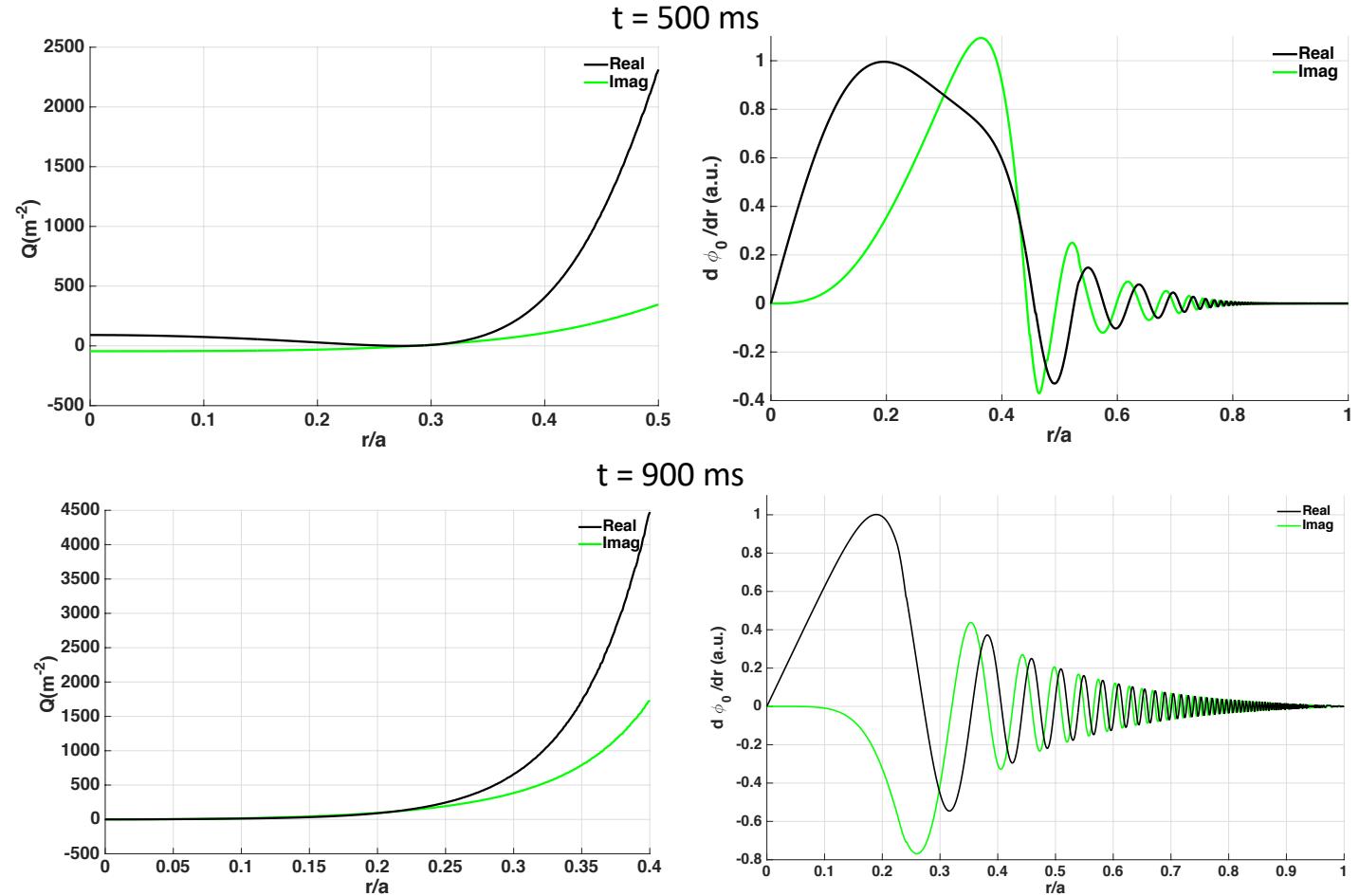
- The local dispersion relation for 500ms and 900ms, injection angle  $\lambda_0 = 0.5$  and  $n_b/n_i \approx 4\%$  for  $r/a < 0.45$ , smoothly decrease for  $r/a > 0.45$ , for 900ms EP is limited to  $r/a < 0.35$ .

- **Local th GAM** is calculated by NOVA.
- The interaction between **EP transit frequency** and **Local th GAM** produces the **Local EGAM**.
- **Maximum OFF-AXIS** in the dispersion relation is found for 500ms, around  $r/a=0.3$ .
- **Maximum ON-AXIS** is found for 900ms.
- Both have positive growth rate  $\gamma > 4\%$ .



# Global E-GAM radial structure

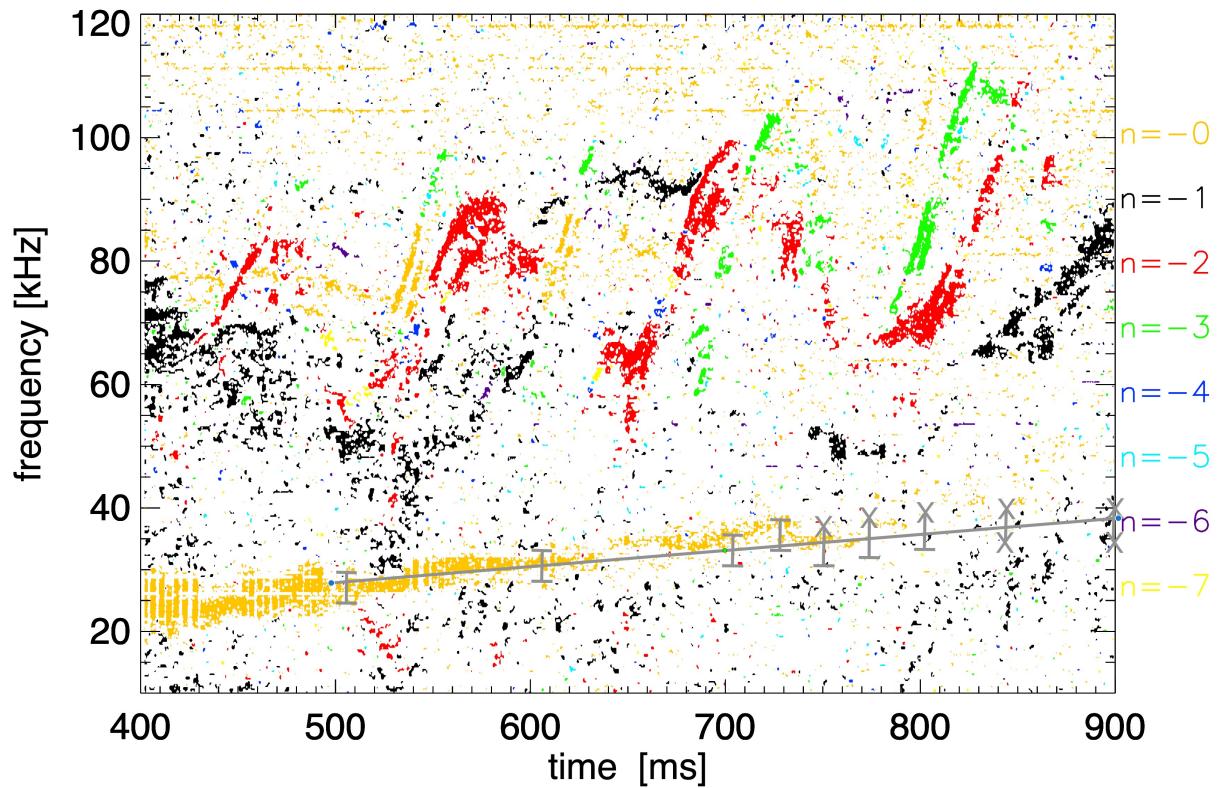
- The global behavior is much more pronounced for 500ms, with Airy function for  $r/a = 0 - 0.5$ , and a outgoing wave is find for  $r/a > 0.5$ .
- For 900ms is found the outgoing wave from  $r/a=0$ .
- In both cases  $\rho^2 k_r^2 \approx 0.6$
- Radial structure similar to obtained by kinetic, numerical [G. Y. Fu, (2008)] and analytically [Z. S. Qu, (2017)] for EGAM.
- MHD numerically [H.L. Berk, 2006] and analytically [Lakhin, 2014], present Dirac for  $d\phi_0/dr$ , in the maximum of dispersion relation,  $M=1,2$  are global.



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# Consistency with N=0 oscillations observed in DIII-D

- NOVA simulation of TRANSP kick model profiles.
- The results match the observed N=0 mode by Mirnov coils on #159243 in DIII-D. The observed frequency in yellow, and the matched points in gray.
- Error bars due to variation of  $\lambda_0 = 0.5 - 0.6$ .
- The frequency increment in time is due to time evolution of temperature and q.
- X means that there is no maximum in the continuum and no eigenmode should be observed at that point.
- Good numerical agreement between theoretical prediction and experimental data.



# Summary

- Slightly reverse  $q$  profile is a source for global EGAM.
- The instability is dominated by injection angle  $\lambda_0$  and absolute value of  $q$ , due to relation of toroidal EP transit frequency and GAM frequency.
- The condition for instability is the presence of OFF-AXIS maximum in the EGAM local dispersion relation.
- Good quantitative agreement with DIII-D N=0 oscillation on #159243 discharge.

# Referencias

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## Profiles used slides 12 and 13

Consistent with TRANSP

$$T_{500}(J) = 1.8(1 - (r/a)^2)^{2.5} \text{ keV}$$

$$T_{900}(J) = 2.0(1 - (r/a)^2)^{2.5} \text{ keV}$$

$$q_{500} = 4.7 - 10(r/a)^2 + 14.8(r/a)^3 - 2(r/a)^4$$

$$q_{900} = 3 - 1.8(r/a) - 4(r/a)^2 + 9.6(r/a)^4$$

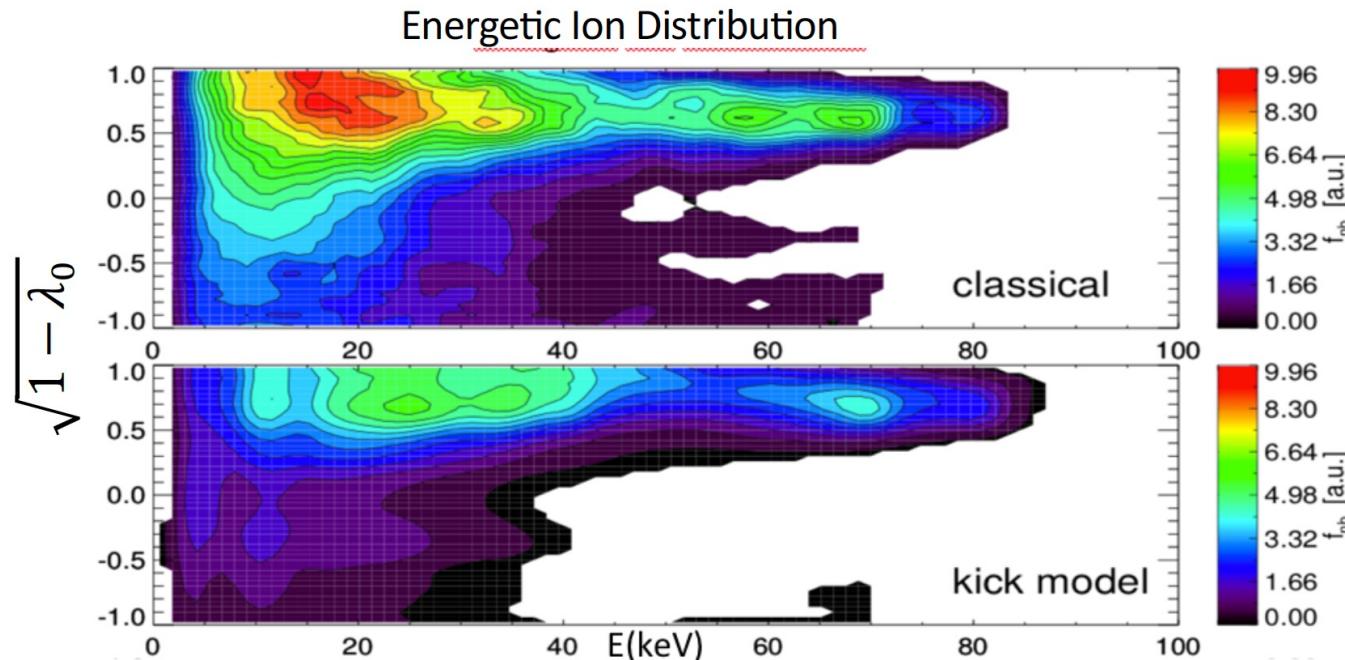
$$\tau_e = 0.86$$

$\sigma_{mhd} = 0.7$  (NOVA gives a profile 0.7 in center 0.75 near edge)

$$R_0 = 1.67m$$

$$a = 0.64$$

# Classical EP profile – TRANSP reconstruction

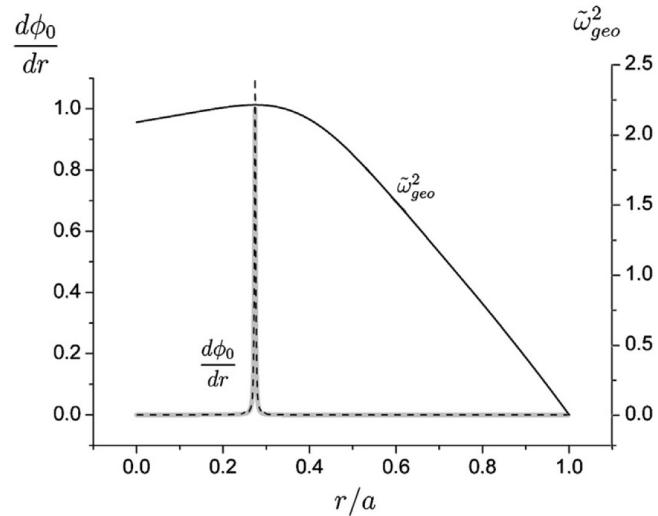


N.N. Gorelenkov et al (2018)

Injection starts with  $\lambda_0 = 0.6$ , and moves to parallel motion as slow down

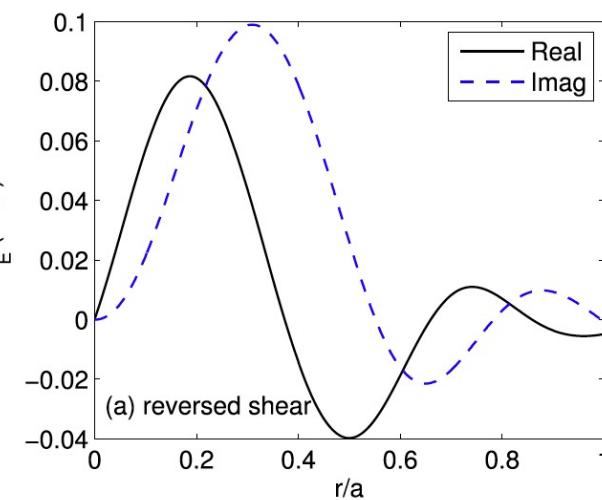
# Eigenmode – previews work comparision

MHD, thermal GAM



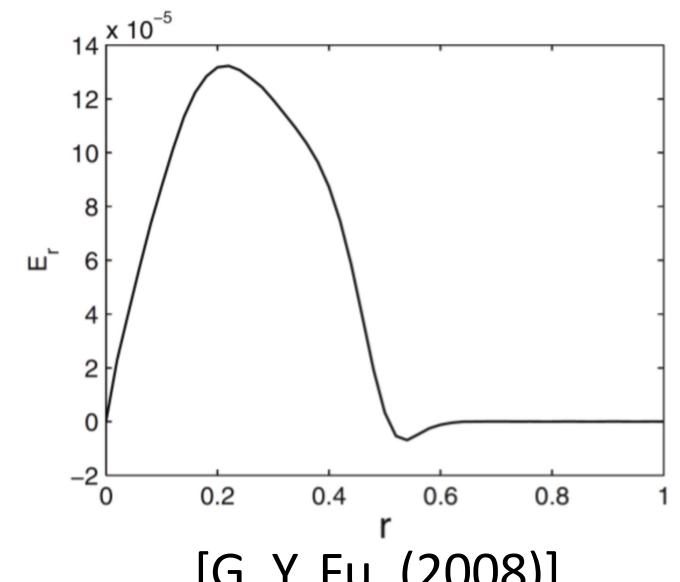
[Lakhin, 2014]

*Kinetic – analytical  
Single energy EP*



[Z. S. Qu, (2017)]

*Kinetic – numerical  
Slowing Down EP*



[G. Y. Fu, (2008)]